

The Conformation of Non-Aromatic Ring Compounds. VII.† Crystal Structure of *trans*-2,5-Dichloro-1,4-dioxane at $-125\text{ }^{\circ}\text{C}$

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trans-2,5-Dichloro-1,4-dioxane crystallizes in the triclinic system: space group $P\bar{1}$, $Z=1$,

$$a=4.573, b=5.491, c=6.665 \text{ \AA}; \alpha=86.25^{\circ}, \beta=104.10^{\circ}, \gamma=106.23^{\circ}$$

at room temperature. Three independent sets of data were used to refine the atomic coordinates:

- (i) Reflexion intensities of the zones [100] and [001], with Mo $K\alpha$ radiation at room temperature.
- (ii) Reflexion intensities of the zones [100] and [001], with Mo $K\alpha$ radiation at $-125\text{ }^{\circ}\text{C}$.
- (iii) Three-dimensional reflexion intensities observed with Cu $K\alpha$ radiation at $-125\text{ }^{\circ}\text{C}$.

The dioxane ring has the chair form with the chlorine atoms in axial positions. Some bond lengths show interesting deviations from the expected values:

$$\text{C}(1)\text{-Cl}=1.845 \pm 0.006, \text{C}(1)\text{-O}=1.388 \pm 0.007, \text{C}(2)\text{-O}=1.428 \pm 0.007,$$

$$\text{C}(1)\text{-C}(2')=1.513 \pm 0.008 \text{ \AA} \text{ at } -125\text{ }^{\circ}\text{C}.$$

1. Introduction

The present study is the first of a series of crystal structure analyses of halogeno-1,4-dioxanes initiated in this laboratory during the past few years. The properties and conformation of this class of substituted heterocyclic molecules have also been investigated in solution by infrared and Raman spectroscopy and by measurements of electric dipole moments (Altona, Romers & Havinga, 1959; Altona, 1963). This research was undertaken in order to determine the preferred conformation or conformations and to compare the results with those from earlier studies of analogous cyclohexane derivatives, both in solution (Kwestroo, Meijer & Havinga, 1954) and in the solid state (Hassel & Vihovde, 1953).

A *cis* and a *trans* form of 2,5-dichlorodioxane are theoretically possible. Only one isomer is known (Bryan, Smedley & Summerbell, 1950). The *trans* configuration has been inferred from the low value (0.60 D) of the dipole moment (Altona, Romers & Havinga, 1959) and is now confirmed by X-ray analysis.

The structure was first refined in projections from molybdenum radiation data obtained at room temperature. Because the carbon-chlorine bond appeared to be unusually long, we decided to improve on the accuracy of the determination by application of Mo $K\alpha$ radiation at low temperature. After completion of this second analysis, digital computing facilities became available and a third refinement was carried out with full three-dimensional Cu $K\alpha$ data, also

collected at $-125\text{ }^{\circ}\text{C}$. These three determinations offer interesting material for comparative purposes.

2. Experimental

trans-2,5-Dichloro-1,4-dioxane, $\text{C}_4\text{H}_6\text{O}_2\text{Cl}_2$, m.p. 122–124 $^{\circ}\text{C}$, was prepared by chlorination of a mixture of 1,4-dioxane and carbon tetrachloride at 10 $^{\circ}\text{C}$ in the presence of a trace of dibenzoyl peroxide (Bryan, Smedley & Summerbell, 1950). The compound crystallizes from carbon tetrachloride in the form of white triclinic laths with principal zone [100] and faces {010} and {001}. Plates with main faces {010} were also found. The crystal faces are rough and the angles α , β and γ were not measured directly with an optical goniometer. The specimens chosen for X-ray diffraction were carefully sealed in dry Lindemann glass capillaries, because the crystals decompose rapidly in the presence of traces of moisture. They were also dipped into liquid air in order to minimize primary extinction effects.

Unit cell dimensions and angles at 20 $^{\circ}\text{C}$ and at $-125\text{ }^{\circ}\text{C}$ were measured from zero-level Weissenberg photographs about [100], [010] and [001], taken with unfiltered copper radiation (λ Cu $K\alpha=1.5418 \text{ \AA}$) and superposed with aluminum powder diffraction lines ($a=4.0489 \text{ \AA}$ at 20 $^{\circ}\text{C}$). No discontinuous structure transition was found between 20 $^{\circ}\text{C}$ and $-125\text{ }^{\circ}\text{C}$. It is interesting to note (Table 1) that angle α^* changes from acute to obtuse on cooling and special care was exercised in indexing the films. The triclinic unit cell contains one molecule. The molecular symmetry is an inversion centre if the space group is $P\bar{1}$, an assumption confirmed by the straightforward refinement.

All intensity records were obtained on non-inte-

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Table 1. *Unit cell dimensions, angles, and densities at 20 °C and at -125 °C*

	20 °C	-125 °C
<i>a</i>	4.573 ± 0.01 Å	4.521 ± 0.007 Å
<i>b</i>	5.491 ± 0.01	5.432 ± 0.01
<i>c</i>	6.665 ± 0.01	6.616 ± 0.01
α	86.25 ± 0.1°	85.73 ± 0.1°
β	104.10 ± 0.15	103.85 ± 0.2
γ	106.23 ± 0.2	106.40 ± 0.25
<i>d</i> (100)	4.268 ± 0.005 Å	4.222 ± 0.003 Å
<i>d</i> (010)	5.272 ± 0.005	5.211 ± 0.005
<i>d</i> (001)	6.464 ± 0.005	6.423 ± 0.005
α^*	89.85 ± 0.05°	90.42 ± 0.05°
β^*	76.40 ± 0.1	76.80 ± 0.1
γ^*	74.20 ± 0.2	74.13 ± 0.2
<i>U</i>	155.8 ± 0.4 Å ³	151.3 ± 0.4 Å ³
<i>D_x</i> (<i>Z</i> =1)	1.676 g.cm ⁻³	1.727 g.cm ⁻³
<i>D_m</i>	1.65	—

grated equi-inclination Weissenberg photographs by the application of the multiple film technique. The films used for exposures with Mo *K* α radiation were interleaved with thin copper foils in order to reduce the intensities on successive films. The low temperature exposures were made according to the method described by Kreuger (1955). Since the exposure times were long (Table 2), an automatic filling system was designed in order to keep the liquid nitrogen in the cooling jacket at a constant level. A full description of the technical details will be given elsewhere (Altona, 1963).

Apart from preliminary Weissenberg photographs taken with copper radiation at room temperature, three sets of data were collected:

1. Reflexion intensities of the zones [100], [001] and [101] with Mo *K* α radiation at room temperature.
2. Reflexion intensities of the zones [100] and [001] with Mo *K* α radiation at -125 °C.
3. Reflexion intensities of the layers 0-3 about [100] and 0-4 about [001] with Cu *K* α radiation at -125 °C.

Some technical details of these sets of exposures are given in Table 2.

Within the Cu *K* α limiting sphere about 690 reflexions are possible; 653 reflexions were recorded, 35 of which were too weak to be observed.

The visually estimated intensities were corrected in the usual way for Lorentz, polarization, and spot-shape factors. The linear absorption coefficient μ (Mo *K* α) = 9.5 cm⁻¹ and μ (Cu *K* α) = 85 cm⁻¹. It follows that the specimens chosen for Mo *K* α exposures (Table 2) were sufficiently small to neglect absorption errors. An absorption correction (cylindrical approximation) was applied to the intensities obtained from the Cu *K* α exposures.

The (*0kl*) and (*h**k*0) observed structure factors of set (3) were put on an approximately absolute scale by comparison with the corresponding calculated structure factors *F_c* of set (2); all higher levels of the former set were then scaled by a thorough cross-correlation procedure. About 360 structure factors *F_o* occurred on the *a*-axis as well as on the *c*-axis photographs. Their mean value *F_m* was compared with the difference *F_m* - *F_o*. The 'agreement factor'

$$R_m = \Sigma |F_m - F_o| / \Sigma F_m = 4.9\%$$

gives an indication of the accuracy of the measurements.

3. Two-dimensional refinements

The simple structure (Fig. 1) became evident after solution of the Patterson projections *P(uv)* and *P(vw)*. With *Z*=1 the asymmetric unit consists of one chlorine atom, one oxygen atom, two carbon and three hydrogen atoms. The numbering of the atoms can be seen in Fig. 3.

Refinement of the structure by means of the room temperature Mo *K* α data (set 1) proceeded by calculation of successive cycles of structure factors, Fourier syntheses and difference maps in projections along [100] and [001]. An anisotropic temperature factor

$$\exp [-(B_1 + B_2 \cos^2 \varphi) (\sin \theta / \lambda)^2]$$

was introduced for the chlorine atom in the *F*(*0kl*) structure factor calculations, the angle between the direction of maximum vibration and the *z* axis being 48 °C. The hydrogen atoms were incorporated into the molecular model by accepting C-H bond distances of 1.03 Å and tetrahedral bond angles. Final agreement indices *R*, taken over all observed reflexions, were:

$$R(0kl) = 5.9\% \quad \text{and} \quad R(hk0) = 8.5\% .$$

Table 2. *Technical data concerning the three sets of photographs taken†*

Set	Radiation	<i>T</i>	Reflexions	Crystal cross-section	Number of observed reflexions	Exposure time	sin θ / λ max
1	Mo <i>K</i> α	20°	<i>0kl</i>	0.20 × 0.35 mm	100	64 hr	0.83
			<i>hk0</i>	0.17 × 0.23	74	180	0.84
2	Mo <i>K</i> α	-125°	<i>0kl</i>	0.30 × 0.30	161	64	1.02
			<i>hk0</i>	0.25 × 0.30	96	92	0.96
3	Cu <i>K</i> α	-125°	<i>0kl</i> - <i>3kl</i>	0.15 × 0.15	618	4-8	0.64
			<i>hk0</i> - <i>hk4</i>	0.15 × 0.15			

† See text.

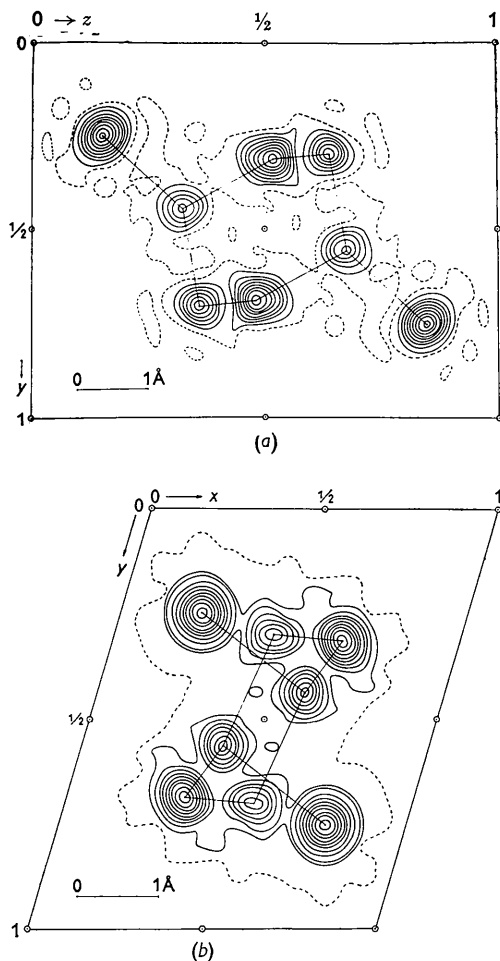


Fig. 1. Electron density projection of *trans*-2,5-dichloro-1,4-dioxane (a) along [100], calculated from the low temperature Mo $K\alpha$ data (set 2); (b) along [001]. Contours are drawn at 2, 4, 6, ... e.Å⁻²; those about the chlorine atom are drawn at 4, 8, 12, ... e.Å⁻². Zero and negative contour lines are broken.

The low-temperature structure (Mo $K\alpha$ data, set 2) was refined in much the same way. Fourier maps (Fig. 1) and a number of difference syntheses were calculated at intervals of 1/120. The direction of maximum vibration of the chlorine atom in projection along [100] was the same as noted previously. All hydrogen atoms were located close to the expected positions in difference maps that included low-angle reflexions only (Jellinek, 1958). Theoretical hydrogen coordinates, however, were used in the final structure factor calculations.

The value of $R(0kl)$ was 8.6% at this stage. Because $F_o < F_c$ for many strong reflexions, a correction for primary extinction was applied (Vand, 1955). The final R indices were: $R(0kl) = 6.3\%$ and $R(hk0) = 6.9\%$. The results of the two analyses will be discussed in § 5 and 6. The observed and calculated structure factors of sets (1) and (2) will be given elsewhere.

The two-dimensional calculations were done with

a desk calculator and Beevers-Lipson strips; for both sets a final cycle of structure factors and difference syntheses was carried out on the electronic XI computer of the University of Leiden, using programs devised by Keulemans (1961). Tables of scattering factors of chlorine (Dawson, 1960), carbon and oxygen (Berghuis, Haanappel, Potters, Loopstra, MacGillavry & Veenendaal, 1955) and hydrogen (McWeeny, 1951) were used throughout this investigation.

4. Three-dimensional refinement

The data of set (3) were used to refine positions, scale-factor, and individual anisotropic temperature parameters by means of least-squares calculations. A block-diagonal approximation to the matrix of the normal equations was applied (Rollett, 1961). The calculations were carried out on the Mercury computer of the University of Manchester, England.

In the first cycle the final atomic positions and the isotropic temperature factors of the two-dimensional refinement with the data of set (2) were inserted. The few non-observed reflexions were given the value $\frac{1}{2}F_{\min}$ where F_{\min} is the locally observable minimum of F_o . The $(F_o - F_c)^2$ values were weighted with a factor:

$$w = 1/[1 + \{(F_o - 10)/5\}^2].$$

The previously assumed hydrogen parameters were included in the calculations, but were not refined.

After four cycles of calculations the shifts in the atomic positions were less than 0.001 Å. Relatively large discrepancies still remained between F_o and F_c of most strong low-order reflexions. Attempts at applying extinction corrections according to the method proposed by Vand (1955) were not successful, presumably because different extinction and residual absorption errors in the data of the two crystals used have been mixed in the three-dimensional analysis.

The observed and calculated structure factors are listed in Table 3. The final agreement index is 8.6%. Exclusion of the non-observed terms and of the nine reflexions most suffering from extinction yields an R value of 7.7%.

Since no positional standard deviations were provided by the least-squares program, several three-dimensional Fourier sections through each atom were calculated at intervals of 1/360. From these maps the central curvatures $\partial^2\rho/\partial x^2$ etc. were obtained. The hydrogen maxima were found to lie close to the expected positions with peak densities 1.0–1.5 Å⁻³, but no new hydrogen parameters were derived.

5. Atomic parameters and accuracy

Table 4 presents the following information, obtained from the three independent sets of data: the final atomic coordinates, the root mean square (r.m.s.) of the nearly isotropic coordinate standard deviations

Table 3. Observed and calculated structure factors (set 3)

Each group of reflexions is preceded by a ● and by the values of h and k common to that group. The remaining lines show l , F_o and F_c ($\times 50$). Non-observed reflexions, inserted as $\frac{1}{2}F_{\min}$, are indicated by X

●	0	-5	-37X	-11	3	622	522	3	-3	3	-315	-322	-6	99	104	0	4	-2		
1	-405	-499	-4	-157	-176	4	675	639	-7	-468	-403	4	248	155	-5	450	415	-4	308	381
2	-229	-269	-3	332	-583	5	-122	-138	-6	193	194	5	535	543	-4	549	544	-3	186	194
3	393	-13	-2	-121	-229	6	-248	-146	-5	648	600	3	571	569	-3	-302	-254	-2	-454	-415
4	-490	-49	-1	121	155	7	-414	-470	-4	239	102	-5	-22X	-4	-2	208	79	-1	-468	-437
5	-212	-101	0	576	610	8	3	3	-3	490	443	-4	-58	-53	-2	-1097	-1132	0	-369	-338
6	737	770	1	221	59	-7	-892	-272	-2	-854	-796	-3	549	580	0	-697	-734	2	307	396
7	120	158	2	99	81	9	81	91	0	-895	-784	-2	108	108	1	382	367	3	276	274
8	268	72	3	-405	-405	-5	603	571	0	-20	-79	-1	-198	-172	2	594	570	3	169	148
-7	387	388	4	-297	-297	4	364	387	2	666	643	0	-333	-345	3	162	177	4	-203	-203
-6	359	359	5	54	65	-3	-85	-84	3	243	225	2	-165	-166	4	153	120	5	4	0
-5	418	417	6	-4	-4	-2	-650	-654	3	616	585	2	-180	-168	5	220	220	6	-121	-126
-4	409	409	7	135	174	-1	-746	-728	4	607	522	1	-74	-72	6	399	399	7	-370	-322
-3	414	414	8	-215	-215	0	-274	-272	5	-245	-200	3	571	569	7	569	569	8	-607	-623
-2	483	483	9	-253	-266	1	796	807	6	-202	-204	4	212	218	8	306	285	9	378	344
-1	211	193	10	-643	-672	2	724	668	7	-103	-120	-3	220	302	9	45X	123	10	329	208
0	2093	2467	11	-263	-272	3	337	330	8	-103	-120	-4	-171	-149	10	-207	-207	11	346	348
1	2220	1552	12	616	603	4	-746	-728	9	-74	-74	-5	-553	-522	11	-553	-522	12	193	175
2	-842	-925	13	532	505	5	-324	-339	10	-212	200	-1	-270	-309	12	-702	-664	13	-261	-234
3	-342	-287	14	-190	-190	6	-346	-349	11	-432	390	0	-153	-128	13	544	468	14	-445	-426
4	-1220	-1392	15	-54	-41	7	126	134	12	-436	-407	2	216	206	0	463	483	15	1	1
5	-198	-162	16	89	905	8	-298	-296	13	298	-296	3	400	399	1	412	412	16	-228	-228
6	442	365	17	121	968	9	270	275	14	-6	-6	4	262	262	2	262	262	17	-122	-122
7	800	790	18	220	236	10	222	205	15	639	580	5	-85	-92	3	-396	-390	18	486	514
8	87	90	19	252	278	11	265	290	16	634	567	-3	-144	-133	4	-558	-539	19	40X	73
9	-99	-124	20	94	126	12	-504	-522	17	571	505	-2	297	279	5	65	71	20	445	402
10	-297	-188	21	-2	-2	13	-212	-212	18	-212	-212	-3	139	139	6	139	139	21	-374	-270
11	198	207	22	-7	-85	14	-100	-120	19	-477	-484	4	-129	-296	7	226	152	22	-256	-231
12	-396	-406	23	-166	-180	15	371	324	20	-58X	-34	5	-157	-168	8	-594	-627	23	3	3
13	-745X	305	24	580	606	16	824	822	21	355	328	6	-328	-373	9	-544	-627	24	4	4
14	293	262	25	637	616	17	-549	-534	22	-1	-1	7	-328	-373	10	222	222	25	5	5
15	1043	1086	26	67	46	18	-351	-345	23	378	388	8	-49	-69	11	535	532	26	30X	56
16	-112	-57	27	202	157	19	-418	-422	24	-726	140	9	58	68	12	227	222	27	-212	-212
17	-265	-232	28	-1220	-1223	20	94	-132	25	54X	72	10	276	360	13	-162	-144	28	-391	-423
18	-148	-160	29	-160	-150	21	350	394	26	-780	-780	11	-162	-162	14	-162	-162	29	-305	-280
19	-257	-147	30	472	441	22	-2	-2	27	-288	-249	12	45X	42	15	-594	-582	30	-27X	-256
20	-49	-10	31	412	401	23	-54	-74	28	-612	-525	13	-603	-556	16	-192	-192	31	280	256
21	904	925	32	391	412	24	-450	-437	29	122	132	14	-186	-206	17	222	222	32	3	3
22	571	568	33	-40X	52	25	-162	-138	30	52	52	15	-40X	-22	18	172	172	33	4	4
23	442	422	34	-256	-304	26	-304	-304	31	528	506	16	32	298	19	32	32	34	5	5
24	-535	-515	35	-207	-288	27	468	449	32	-827	-772	17	94	153	20	-128	-128	35	6	6
25	-252	-223	36	-8	-8	28	567	530	33	-540	-525	18	-4	-4	21	-124	-124	36	7	7
26	192	192	37	-193	-213	29	99	99	34	737	691	19	-16	-16	22	-346	-346	37	8	8
27	883	241	38	-6	-6	30	-594	-606	35	-152	-152	20	316X	316	23	265	265	38	9	9
28	265	245	39	442	444	31	-139	-135	36	308	299	21	504	504	24	364	364	39	10	10
29	499	545	40	684	641	32	-474	-462	37	283	226	22	220	216	25	-207	-207	40	11	11
30	-197	-162	41	-292	-264	33	333	333	38	-1	-1	23	256	234	26	-897	-877	41	12	12
31	-585	-588	42	-508	-485	34	387	435	39	-184	-292	24	-10	-12	27	-472	-472	42	13	13
32	-729	-678	43	-1506	-1742	35	-72	-86	40	-378	-344	25	-657	-657	28	36X	33	43	14	14
33	352	324	44	666	613	36	67	95	41	-450	-391	26	49X	15	3	320	320	44	-108	-199
34	52	52	45	52	52	37	58	52	42	418	374	27	228	222	29	297	334	45	-162	-158
35	2259	1978	46	1024	1166	38	-128	-128	43	123	1150	28	328	328	30	4	4	46	220	236
36	-49	-50	47	-175	-120	39	-319	-286	44	63	34	29	-139	-122	31	94	92	47	346	364
37	-445	-432	48	-126	-128	40	-274	-264	45	-882	-882	30	-225	-169	32	-130	-122	48	-279	-296
38	-477	-480	49	-486	-505	41	-325	-315	46	944	-874	31	944	-874	33	-220	-220	49	-279	-296
39	-112	-125	50	-267	-272	42	40X	40X	47	122	99	32	247	260	34	265	267	50	-279	-296
40	388	277	51	67	67	43	391	437	48	940	901	33	391	366	35	265	266	51	3	3
41	139	155	52	433	383	44	243	165	49	126	205	34	103	120	36	342	346	52	4	4
42	-427	-422	53	477	472	45	6	6	50	274	284	35	275	149	37	372	346	53	31X	29
43	-459	-459	54	-265	-319	46	-247	-274	51	-378	-329	36	-1079	-1122	38	-1122	-1122	54	-477	-478
44	45X	0	55	-670	-670	47	-123	-123	52	-279	-377	37	-40X	-22	39	-126	-126	55	-216	-202
45	36X	40	56	-640	-640	48	261	279	53	-272	-205	38	585	521	40	-280	-281	56	-2	-2
46	728	692	57	-791	-822	49	166	213	54	-225	-214	39	160	392	41	-5	-5	57	288	340
47	612	567	58	520	675	50	-77	-77	55	-162	-152	40	234	212	42	418	483	58	274	263
48	-528	-420	59	1264	1166	51	-166	-278	56	890	893	41	-622	-666	43	-63	-73	59	-222	-222
49	-412	-412	60	769	861	52	-6	-6	57	157	127	42	-162	-172	44	-215	-227	60	-225	-220

Table 4. *Final atomic coordinates in fractions of the cell edges*

The root mean square standard deviations (Å), and orthogonal atomic coordinates (Å) from the three sets of data; the hydrogen positions are theoretical (see text)

Atom	Set	<i>y/a</i>	<i>x/b</i>	<i>z/c</i>	σ	<i>X</i>	<i>Y</i>	<i>Z</i>
Cl	1	0.2290	0.2495	0.1545	0.0030 Å	0.7638 Å	1.1447 Å	0.9987 Å
	2	0.22950	0.24900	0.15400	0.0018	0.7634	1.1355	0.9892
	3	0.23085	0.24881	0.15462	0.0013	0.7684	1.1330	0.9932
O	1	0.6630	0.3155	0.5155	0.0070	2.105	1.110	3.332
	2	0.6640	0.3135	0.5140	0.0045	2.106	1.109	3.302
	3	0.6648	0.3133	0.5129	0.0037	2.103	1.109	3.295
C(1)	1	0.5940	0.4440	0.3260	0.0110	2.098	1.821	2.107
	2	0.5940	0.4450	0.3250	0.0064	2.087	1.819	2.088
	3	0.5936	0.4428	0.3267	0.0056	2.082	1.808	2.098
C(2)	1	0.4540	0.3000	0.6400	0.0100	0.992	1.346	4.137
	2	0.4490	0.2970	0.6435	0.0056	0.978	1.357	4.133
	3	0.4467	0.2969	0.6411	0.0055	0.972	1.358	4.118
H(1)	calc.	0.770	0.470	0.254	—	2.96	1.70	1.63
H(2)	calc.	0.238	0.178	0.578	—	0.16	0.95	3.71
H(3)	calc.	0.548	0.227	0.786	—	1.19	0.92	5.05

Table 5. *Temperature factors (Å²) from the three sets of data*

Atom	Set (1)		Set (2)		Set (3)					
	<i>B</i> [100]	<i>B</i> [001]	<i>B</i> [100]	<i>B</i> [001]	<i>B</i> ₁₁	<i>B</i> ₂₂	<i>B</i> ₃₃	<i>B</i> ₁₂	<i>B</i> ₁₃	<i>B</i> ₂₃
Cl	2.11*	3.85	1.06*	1.33	1.99	2.11	1.75	0.57	0.15	−0.50
O	2.93	3.55	1.13	1.06	1.73	2.27	2.26	0.92	0.65	−0.13
C(1)	2.68	3.81	1.15	1.08	1.76	2.49	1.48	0.68	0.21	−0.29
C(2)	2.75	3.52	1.13	1.26	2.16	2.15	2.10	0.65	0.77	−0.20
H(1), (2), (3)	3.00	3.80	1.10	1.10	isotropic <i>B</i> = 1.10					

* See § 3 for the thermal-anisotropy correction.

anisotropic temperature factors b_{ij} , defined by $\exp -[b_{11}h^2 + 2b_{12}hk + b_{22}k^2 + 2b_{13}hl + 2b_{23}kl + b_{33}l^2]$ were converted into Å² units by means of the relation:

$$B_{ij} = 4b_{ij}/(a_i^* a_j^*) .$$

Although the photographs of sets (2) and (3) have been taken at approximately the same temperature, the temperature factors obtained from the copper radiation data are larger. Probably a number of strong reflexions near the upper edge of the Cu $K\alpha$ films have been underestimated, whereas on the Mo $K\alpha$ films the weak high-order spots may have been assigned too large a value. Whatever the cause, temperature factors obtained from visually estimated intensities clearly have only qualitative significance. In order to check the directions of the thermal motion, the principal axis parameters of the vibration ellipsoid

of the chlorine atom were calculated from the b_{ij} values of set (3) (Busing & Levy, 1958). The directions and magnitudes are:

$$\begin{aligned} U_s &= -0.6208\mathbf{a} + 1.2814\mathbf{b} + \mathbf{c}, & B_s &= 0.55 \text{ \AA}^2 \\ U_m &= 4.9383\mathbf{a} + 0.4421\mathbf{b} + \mathbf{c}, & B_m &= 1.88 \text{ \AA}^2 \\ U_l &= -1.3270\mathbf{a} - 1.6741\mathbf{b} + \mathbf{c}, & B_l &= 2.98 \text{ \AA}^2 \end{aligned}$$

where \mathbf{a} , \mathbf{b} and \mathbf{c} are basic unit cell vectors defined in Table 1.

The angle between the maximum displacement vector and the carbon-chlorine bond is 108.9°. In the projection along [100] the angle between the direction of maximum vibration and \mathbf{c} is 46.5°; this figure agrees with that found (48°) in both two-dimensional analyses.

The observed peak electron densities ρ_0 and the central curvatures A are given in Table 6. These

Table 6. *Observed electron densities and central curvatures from the three sets of measurements*

Atom	e. Å ⁻²		e. Å ⁻³	e. Å ⁻⁴		e. Å ⁻⁵		
	ρ_1	ρ_2	ρ_3	− A_1	− A_2	− A_{3hh}	− A_{3kk}	− A_{3ll}
Cl	28.0 ± 0.2	46.0 ± 0.3	44.0 ± 0.1	520	930	432	460	446
O	12.2 ± 0.2	19.9 ± 0.3	16.4 ± 0.1	180	367	161	154	152
C(1)	8.8 ± 0.2	14.6 ± 0.3	11.2 ± 0.1	116	252	109	99	104
C(2)	10.2 ± 0.2	15.4 ± 0.3	11.4 ± 0.1	138	307	113	103	102
H	0.6 ± 0.2	1.1 ± 0.3	1.4 ± 0.1					

values generally agree with those expected on theoretical grounds (Atoji, 1957). It is well known that, because of series termination effects, lowering of the temperature has a smaller influence on the central curvatures obtained from Cu $K\alpha$ data than on the curvatures from Mo $K\alpha$ data. Our work illustrates that, in some favourable cases at least, the standard errors of a two-dimensional analysis with molybdenum radiation at low temperature are comparable to those of a three-dimensional refinement with Cu $K\alpha$ data at the same temperature.

6. Results and discussion

(a) Geometry of the molecule

The ring in *trans*-2,5-dichloro-1,4-dioxane has the normal chair form and the chlorine atoms occupy axial positions (*aa*). From infrared and Raman spectroscopic data the conclusion was drawn that the same conformation predominates in solution, regardless of the solvent (Altona, Romers & Havinga, 1959; Altona, 1963). In fact, no indications were found for the existence of other conformations, such as the diequatorial one (*ee*). This behaviour is in contrast to that of *trans*-1,4-dichlorocyclohexane, a compound which occurs in solution as an equilibrium mixture containing both the (*ee*) and the (*aa*) chair forms (Kozima, Sakashita & Maeda, 1954), while in the solid state only the (*ee*) conformation has been found (Hassel & Vihovde, 1953). As a rule, the halogen atoms in dihalogeno-dioxanes occupy axial positions; the (*aa*) forms were also found in *trans*-2,3-dichlorodioxane and in *trans*-2,3-dibromodioxane, for which X-ray structure determinations have recently been accomplished (Altona & Romers, 1963; Altona, Knobler & Romers, 1963).

Table 7. Bond distances and bond angles in *trans*-2,5-dichlorodioxane as determined from three independent sets of measurements

Atoms	Distances (Å)		
	Set (1)	Set (2)	Set (3)
C(1)-Cl	1.862 ± 0.011	1.852 ± 0.007	1.845 ± 0.006
C(1)-O	1.417 ± 0.013	1.407 ± 0.008	1.388 ± 0.007
C(2)-O	1.393 ± 0.013	1.423 ± 0.008	1.428 ± 0.007
C(1)-C(2')	1.521 ± 0.015	1.504 ± 0.009	1.513 ± 0.008
	Angles (°)		
C(2')-C(1)-Cl	108.4 ± 0.8	109.2 ± 0.5	109.6 ± 0.4
O-C(1)-Cl	109.5 ± 0.8	109.6 ± 0.5	110.2 ± 0.4
C(1)-O-C(2)	114.2 ± 0.8	113.9 ± 0.6	113.1 ± 0.5
O-C(1)-C(2')	111.4 ± 0.9	112.4 ± 0.6	112.4 ± 0.5
O-C(2)-C(1')	112.8 ± 0.9	111.9 ± 0.6	111.7 ± 0.5

The bond distances and angles, calculated from the parameters in Table 4, are given in Table 7, together with the estimated standard deviations (Cruickshank & Robertson, 1953; Darlow, 1960) in which allowance has been made for the errors in the unit-cell dimen-

sions. No corrections for the effect of molecular librations were applied, mainly because of the uncertainty in the magnitudes of the vibration parameters.

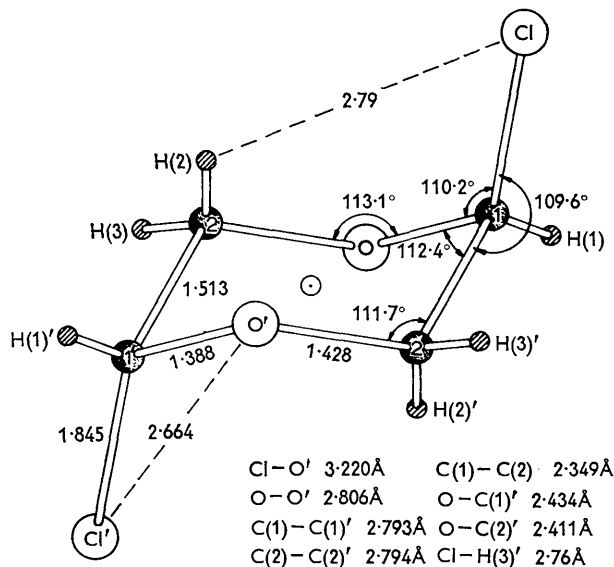


Fig. 2. Schematic drawing of the molecule of *trans*-2,5-dichloro-1,4-dioxane showing some intramolecular distances and angles.

Agreement among the three independent analyses is again gratifying; none of the differences in bond distances between each set is found to be significant. In the following discussion the distances and angles from set (3), also shown in Fig. 2, will be used, as these are considered to be the most reliable values.

The valency angles are normal: those in the ring are 2-3° larger than the tetrahedral values, in accordance with many data on aliphatic ring and open chain compounds; those about the C-Cl bond do not differ significantly from 109.5°. Interesting bond distances are found in this molecule. The C-O bonds are not equivalent: C(1)-O (1.388 Å) is 0.040 Å shorter than C(2)-O (1.428 Å). The latter bond distance does not deviate significantly from the C-O distance found in dimethyl ether (1.416 Å, Kimura & Kubo, 1959; 1.417 Å, Kasai & Myers, 1959), but C(1)-O is appreciably shorter. The C-C bond distance is comparable to that reported for ethyl chloride (1.520 ± 0.003 Å; Schwendeman & Jacobs, 1962); both are perhaps shorter than the C-C distance in normal paraffins (1.533 Å; Bartell, 1959).

The C-Cl bond length is surprisingly large (1.845 ± 0.006 Å). An average C-Cl bond distance of 1.767 Å is given by Sutton (1958). However, this figure includes distances in many compounds that contain two or three chlorine atoms attached to the same carbon atom. Few accurate C-Cl bond distances have been reported for saturated aliphatic compounds of the types -CHCl- and -CH₂Cl. The available results

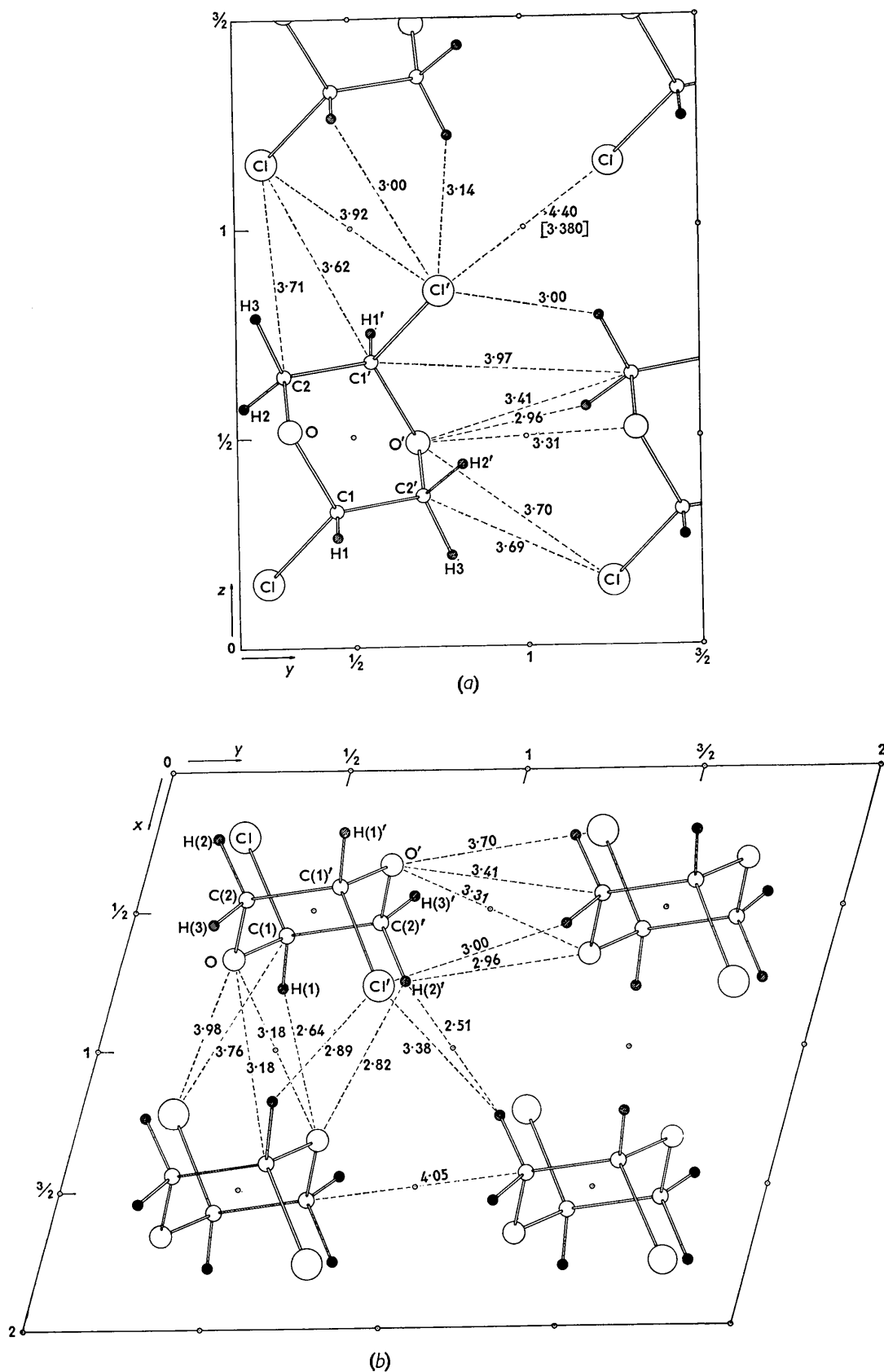


Fig. 3. (a) Orthographic projection along [100] of a layer of unit cells. Some intermolecular distances are shown. Distance between brackets refers to molecule at $x+1, y+1, z+1$. (b) The same along [001].

seem to indicate that the 'normal' aliphatic C-Cl bond has a length of about 1.78-1.79 Å. Examples are: methyl chloride (1.784 Å; Bartell & Brockway, 1955; 1.7815 Å, Miller, Aamodt, Dousmanis, Townes & Kraitchman, 1952) and ethyl chloride (1.788 Å; Schwendeman & Jacobs, 1962). Some examples of longer C-Cl bonds are also known, e.g. in 1-chloro-2-butyne (1.798 Å; Laurie & Lide, 1959) in allyl chloride (1.82 Å; Bowen, Gilchrist & Sutton, 1955), and in acetyl chloride (1.82 Å; Allen & Sutton, 1951). A C-Cl bond distance of 1.87 Å has been reported for chloromethyl methyl ether (Akishin, Vilkov & Sokolova, 1960), but a recent redetermination (Planje, Toneman & Dallinga, 1963) gave the value of 1.81 Å. We also found long C-Cl bonds in *trans*-2,3-dichlorodioxane (mean value of two chemically equivalent bonds 1.835 Å) and in *cis*-2,3-dichlorodioxane where the axial C-Cl bond has a length of 1.819 ± 0.008 Å.

Allen & Sutton (1951) observe in their discussion of the structure of the acetyl halides that all long bonds involving halogen atoms are unusually reactive in hydrolysis and dissociate easily. The halogen atoms in monochlorodioxane and in dichloro- and dibromodioxane, like those in open chain α -halo-ethers, are indeed highly reactive in hydrolysis reactions. Furthermore, the above authors conclude from dipole moment data that the increase in bond length is not correlated with a change in ionic character in the bonds. This conclusion is supported by the value of the dipole moment of monochlorodioxane, which indicates that a description in terms of contributions of mesomeric structures such as



does not seem adequate to account quantitatively for the considerable changes in the bond distances.

A satisfactory explanation for the phenomena cannot be given in the present state of our knowledge, but some points are worth noting. First, there are indications that the presence of lone electron pairs on an atom C in a compound A-B-C lengthens the bond A-B, particularly if the ion A⁻ is a weak base (cf. Bent, 1962). Secondly, the equatorial C-Cl bond in *cis*-2,3-dichlorodioxane is 'normal' in length (1.782 ± 0.008 Å). Therefore, the spatial arrangement of atoms and lone pairs may well be important. It is not known with certainty which orbitals are used by the lone pair electrons of an ether oxygen atom, but according to the available evidence (Dickens & Linnett, 1957; Gillespie & Nyholm, 1957) two lone pairs are probably equivalent. Each oxygen atom of the dioxane ring then carries one equatorial and one axial lone pair; the axial lone pair and the oxygen, carbon and the axial, but not the equatorial chlorine atoms will have a planar configuration. It is known that planar configurations of atoms play an important role in the stereochemistry of elimination reactions. It now appears that coplanarity involving a lone pair at the one end of the system and a chlorine

atom at the other end may produce changes in the bond distances between the atoms in the plane. A more extensive discussion of these problems will be given elsewhere (Altona, 1963).

(b) Molecular environment

All intermolecular distances up to 4.4 Å have been calculated. Some of these are shown in Fig. 3. A fairly short Cl...Cl distance of 3.38 Å is found between two chlorine atoms in approximately 'head-on' positions: Cl(X, Y, Z) and Cl'(X-1, Y-1, Z-1). The angle C(1)-Cl...Cl' is 157.1°. Examples of comparable Cl...Cl separations are given by Fischmann, MacGillavry & Romers (1961) and by Hermodsson (1960). The chlorine atoms are coordinated by six hydrogen atoms, two belonging to the 'original' molecule (H2 and H3, Fig. 3) and four others at distances varying between 2.89 Å and 3.14 Å. The shortest distances among the remaining atoms are: H...H=2.51 Å; O...O=3.18 Å; H...O=2.64 Å; C...O=3.18 Å; Cl...C=3.62 Å. The packing is seen to be relatively dense.

The intra- and intermolecular atomic distances and angles were calculated on the X1 computer by means of an exhaustive program, written in ALGOL 60 by Mr H. L. Jonkers with the collaboration of Dr A. Ollongren and one of us (C. A.) (Program A 200, Centraal Reken-Instituut, The University, Leiden).

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The Conformation of Non-Aromatic Ring Compounds. VIII.* The Crystal Structure of *cis*-2,3-Dichloro-1,4-dioxane at -140°C

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cis-2,3-Dichloro-1,4-dioxane crystallizes in the orthorhombic system, space group $P2_12_12_1$, $Z = 4$, and $a = 4.463$, $b = 10.67$, $c = 13.13$ Å at -140°C .

The structure was refined in three projections by means of data obtained at about -140°C with molybdenum radiation.

The dioxane ring has the chair form. The two carbon-chlorine bonds are unequal in length: C(1)-Cl(1) (axial) = 1.819 ± 0.009 Å, C(2)-Cl(2) (equatorial) = 1.781 ± 0.007 Å.

Introduction

It seemed worthwhile to determine the structure of *cis*-2,3-dichloro-1,4-dioxane by X-ray methods, since this compound is the only known dihalogeno-1,4-dioxane in which an equatorial carbon-halogen bond was expected to be present (Altona, Romers & Havinga, 1959).

Experimental

cis-2,3-Dichloro-1,4-dioxane was prepared by chlorination of a refluxing mixture of 1,4-dioxane and carbon

tetrachloride (Summerbell & Lunk, 1957). The compound crystallizes from carbon tetrachloride as colourless needles, elongated along the a axis. The crystals are soft and tend to disintegrate when even slight pressure is applied in directions perpendicular to the needle axis. After several trials some good crystals were obtained that could be cut along [010] and [001]. Their cross sections were roughly hexagonal with faces {001} and {011} about equally developed. A central capillary hole with a radius of approximately 0.02 mm ran along the full length of all these needles.

The dimensions of the unit cell were determined from zero-layer Weissenberg photographs taken with unfiltered Cu $K\alpha$ radiation ($\lambda = 1.5418$ Å) about [100],

* Part VII: Altona, Knobler & Romers (1963); for earlier parts see van Dort & Havinga (1963).